

Lattice Gauge Theory and High Performance Computing

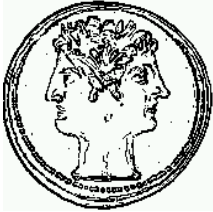
Karl Jansen



- **Introduction**
- **Static potential**
 - why quarks are not free
 - what happens when they try to escape
- **How it works**
discretizing, simulating and back to the continuum
- **More to compute**
 $\alpha_s, m_{\text{quark}}, \langle x \rangle, \dots$
- **Why we are not happy**
 - chiral symmetry
 - Algorithms and machines
- **Outlooks and Needs**

Why Lattice Gauge Theory had to be invented

→ QuantumChromoDynamics

asymptotic freedom		confinement
distances $\ll 1\text{fm}$		distances $\gtrsim 1\text{fm}$
world of quarks and gluons		world of hadrons and glue balls
perturbative description		non-perturbative methods

non-perturbative methods: Lattice
(in combination with e.g. chiral perturbation theory)

- conquering the challenge K.G. Wilson 1974
- demonstrating practicability M. Creutz 1980

Spreading the news around ...

using lattice methods in many areas
different from QCD

- rigorous mathematical definition of field theories
e.g. QCD, chiral gauge theories
- Electroweak physics
e.g. Higgs mass bounds, strength of finite temperature phase transition
- spin models
e.g. precise critical exponents
- Supersymmetry
e.g. phase diagram, mass spectrum
- Quantum gravity
e.g. matrix models, construction of quantum gravity

What Lattice Gauge Theory can provide

- test of theory and validity of analytical methods
 - overlap with
 - *perturbation theory*
 - *chiral perturbation theory*
 - *large- N expansion*
 - *sum rules*
- precise values of many physical observables
 - overlap with
 - *phenomenology*
 - *experiment*
- understanding of strong interaction
 - *semi-classical picture*
 - *monopoles, instantons*

Example: electroweak phase transition

→ test of perturbation theory

exciting possibility: *baryon-asymmetry* of the universe is generated in an early stage of the universe at the *electroweak phase transition* at $T_c \approx 250\text{GeV}$

Condition Sakharov;Kuzmin,Rubakov,Shaposhnikov

- rate of baryon generation \neq rate of baryon annihilation
- out of equilibrium phenomena
- strong enough *first order* phase transition

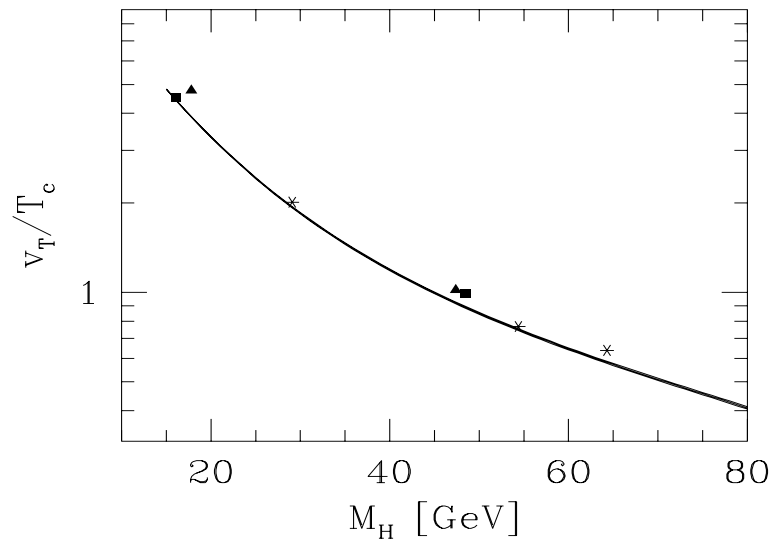
$\frac{v_T}{T_c} > 1$	jump of order parameter v_T large enough
v_T	<i>Higgs vacuum expectation value</i>
T_c	<i>critical temperature</i>

electroweak physics \Rightarrow use perturbation theory

Buchmüller, Fodor, Hebecker

uncertainty in perturbation theory triggered numerical lattice simulations of the electroweak sector (SU(2)-Higgs model)

- 4-dimensional simulations at finite temperature
Fodor, Hein, Jansen, Jaster, Montvay
- 3-dimensional effective field theory simulations
Kajantie, Laine, Shaposhnikov, Rummukainen

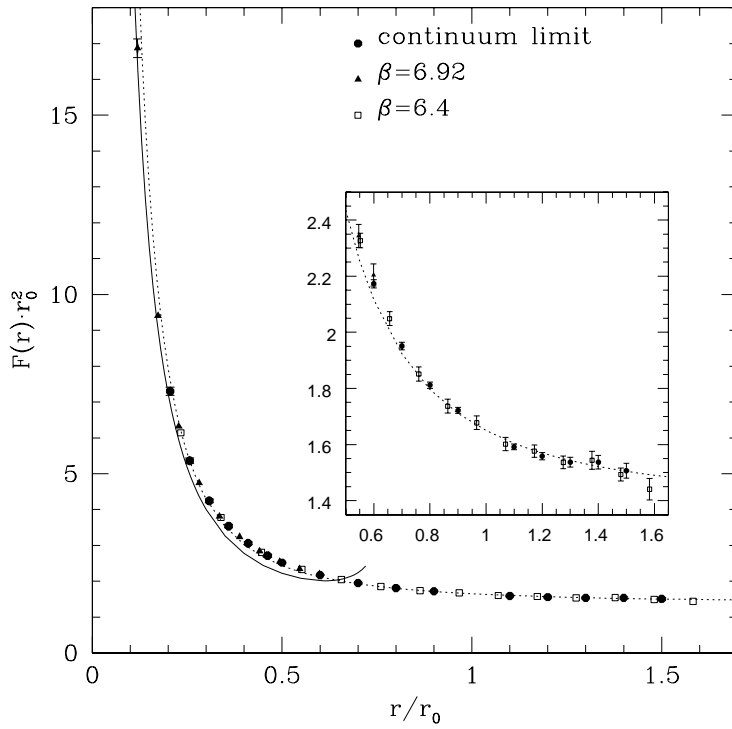


Other examples

- $g - 2$: contribution of light by light scattering
- parton distributions at small momentum transfer
- matrix elements for (indirect) CP-violation
- Kaon decay matrix elements
- \vdots

Test of the pieces

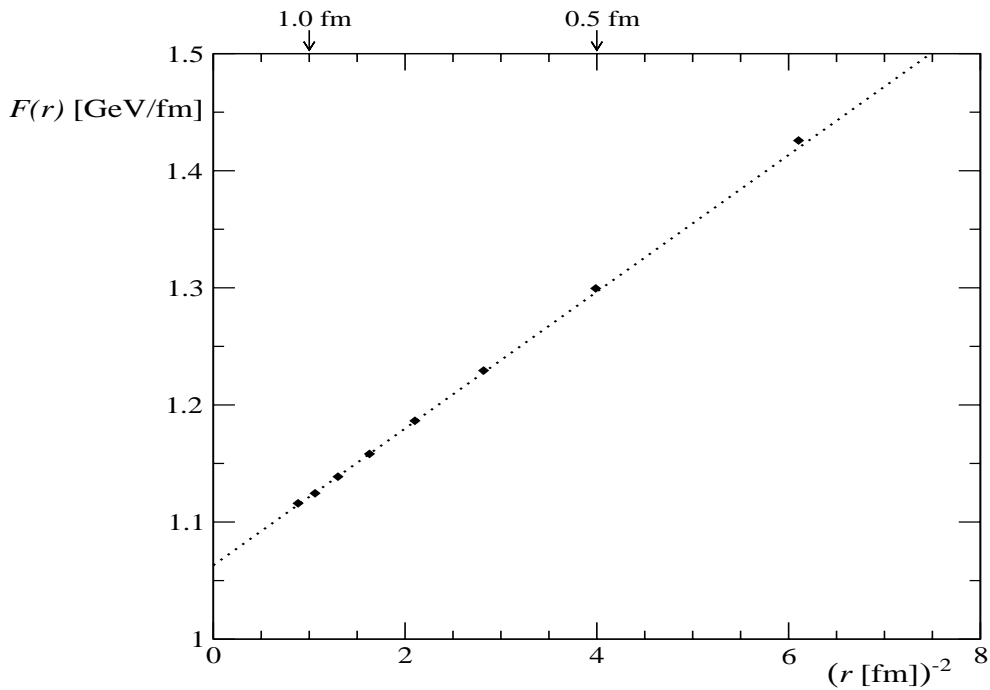
- short to intermediate ranges (Necco and Sommer)



$$r_0 \approx 0.5 \text{ fm}$$

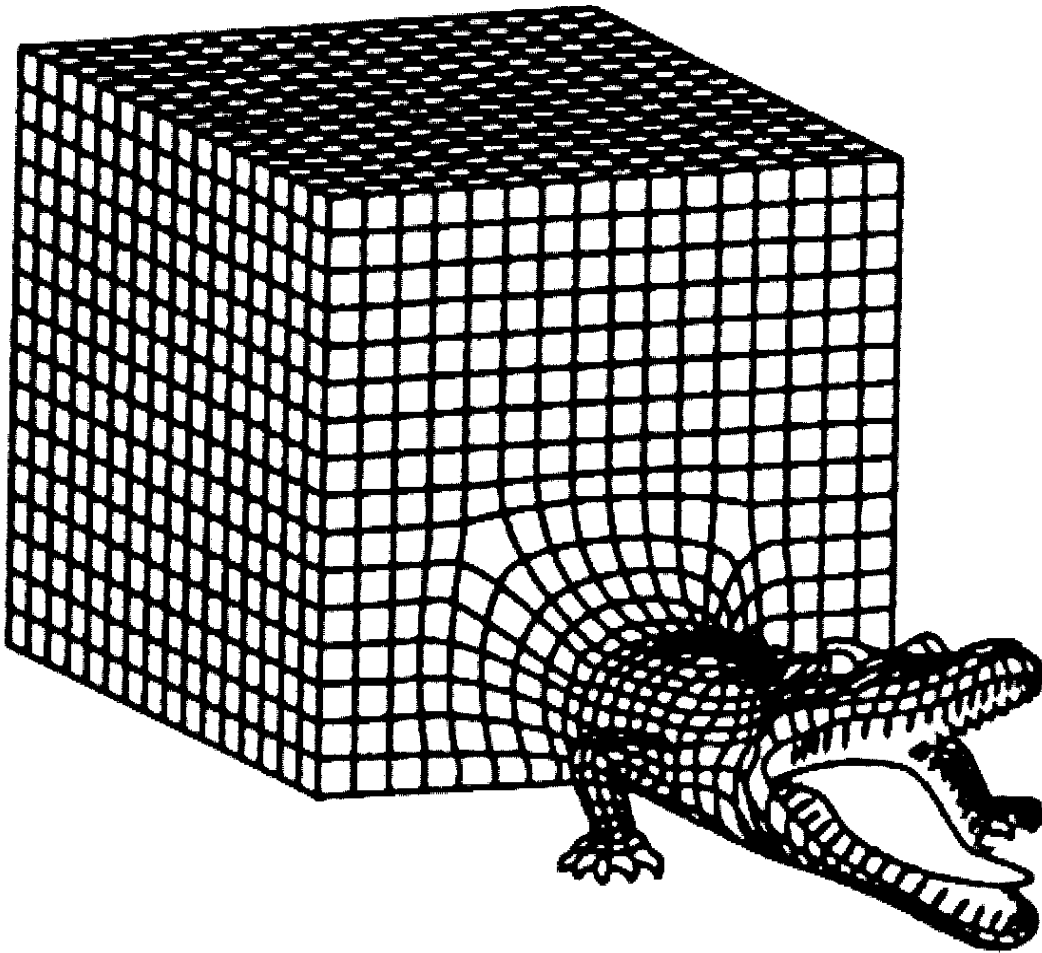
$$F(r) = \alpha \cdot 1/r^2 + \sigma$$

- long range (Lüscher and Weisz)



$$F(r) = \pi/12 \cdot 1/r^2 + \sigma$$

Dangerous lattice Animals



- discretization errors
- finite volume effects

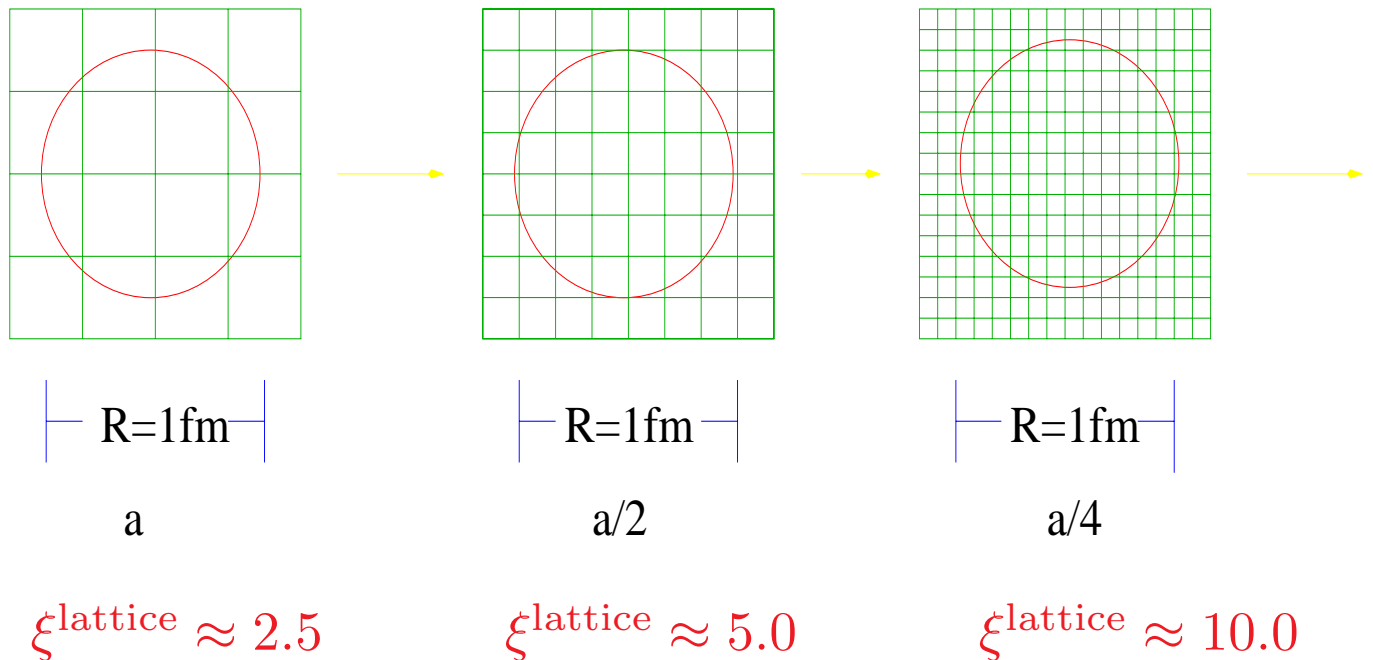
A look at the continuum limit

the general idea of the continuum limit:

we keep fixed values of physical quantities such as a particle mass $m^{\text{phys}} = m^{\text{lattice}}/a$

\Rightarrow for $a \rightarrow 0 \Rightarrow m^{\text{lattice}} \rightarrow 0$

since $m^{\text{lattice}} = 1/\xi^{\text{lattice}}$ in the continuum limit
the *lattice correlation length diverges*



a fixed physical size of a particle (e.g. the proton) receives a finer and finer resolution as $\xi^{\text{lattice}} \rightarrow \infty$

The continuum limit

fixed *physical* length $L = Na = 1\text{fm}$ means

$$a = 0.1\text{fm} \Rightarrow N = 10$$

$$a = 0.05\text{fm} \Rightarrow N = 20$$

$$a = 0.01\text{fm} \Rightarrow N = 100$$

number of lattice points: N^4

easily running out of **computationtime** and **memory**

solutions (?)

- keep $a \gg 0 \Rightarrow$ lattice artefacts
- keep $L < 1\text{fm} \Rightarrow$ finite size effects

modern approach through theoretical advances

- accelerate continuum limit: **improvement programme**
- do not be afraid of finite size effects: **make use of them**

Acceleration to the continuum limit

(old) standard lattice action of QCD is

$$S_{\text{old}} = \underbrace{S_G}_{O(a^2)} + \underbrace{S_{\text{wilson}}}_{O(a)}$$

⇒ expectation values of physical observables

$$\langle O \rangle = \langle O \rangle_{\text{cont}} + O(a)$$

employing all lattice symmetries, equations of motions

⇒ only one more term in $O(a)$ possible

(improved) standard lattice action

Skeikoleslami and Wohlert

$$S_{\text{new}} = S_{\text{old}} + \underbrace{S_{\text{sw}}}_{O(a)}$$

$$S_{\text{sw}} = a^5 \sum_x c_{\text{sw}} \bar{\psi}(x) \frac{i}{4} \hat{F}_{\mu\nu}(x) \Psi(x)$$

with c_{sw} a *tunable* parameter

⇒ compute non-perturbatively c_{sw} such that $O(a)$ cancel

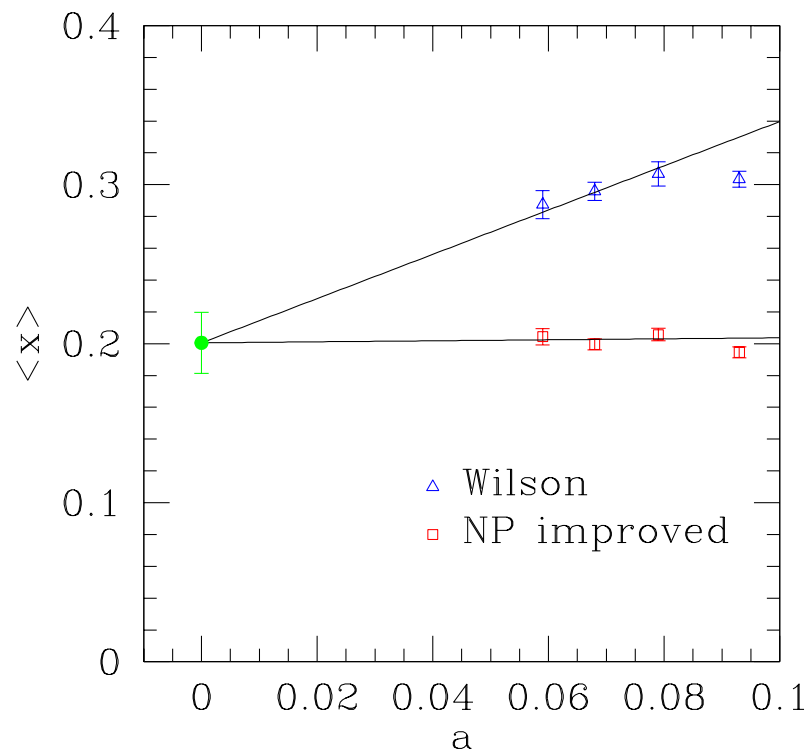
⇒ (nota bene: if also the operator is improved)

$$\langle O \rangle = \langle O \rangle_{\text{cont}} + O(a^2)$$

successful *Symanzik improvement programme* of the



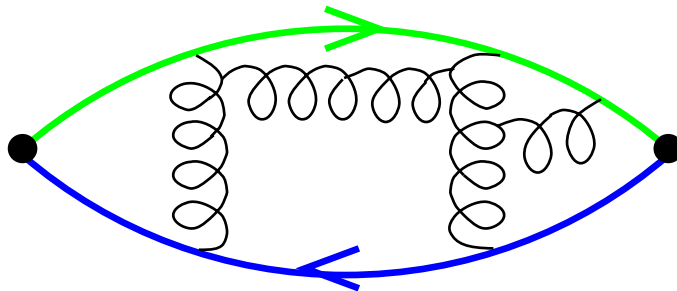
Example of physical quantity: *average momentum* $\langle x \rangle$
of non-singlet, twist-2 operator in a pion



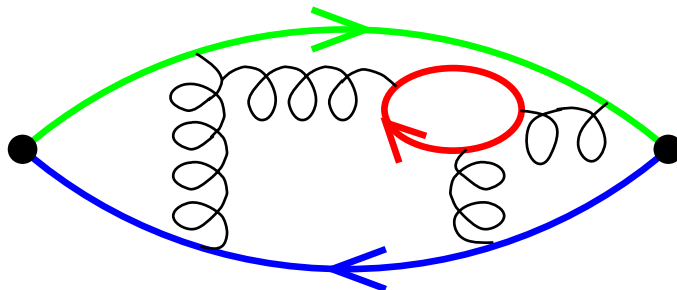
Simulations so far mostly done in

The quenched approximation

- truncation of the theory
- simulations much cheaper
- surprisingly close to experimental values
O(20%) deviation for many quantities



(A) Quenched QCD: no internal quark loops



(B) full QCD

mass spectrum

one of the major goals of lattice QCD

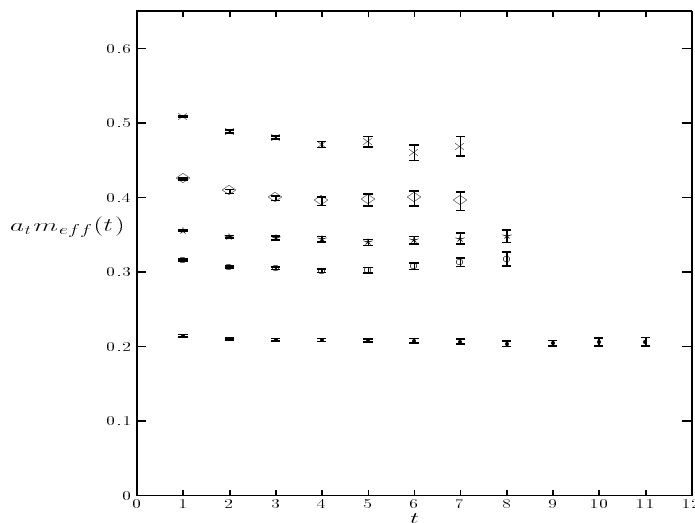
→ compute hadron masses from first principles

define appropriate operators at zero momentum

$$O(t) = \sum_{\mathbf{x}} O(\mathbf{x}, t)$$

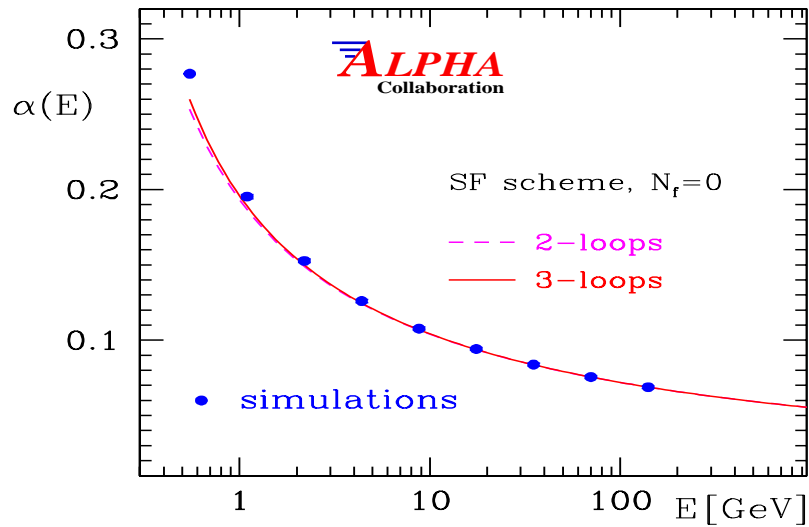
$$\begin{aligned} \langle O(0)O(t) \rangle &= \frac{1}{Z} \sum_n \langle 0|O(0)e^{-\mathbb{H}t}|n\rangle \langle n|O(0)|0\rangle \\ &= \frac{1}{Z} \sum_n |\langle 0|O(0)|n\rangle|^2 e^{-(E_n - E_0)t} \end{aligned}$$

→ effective masses $m_{\text{eff}}(t) = -\ln \frac{\Gamma(t+1)}{\Gamma(t)}$

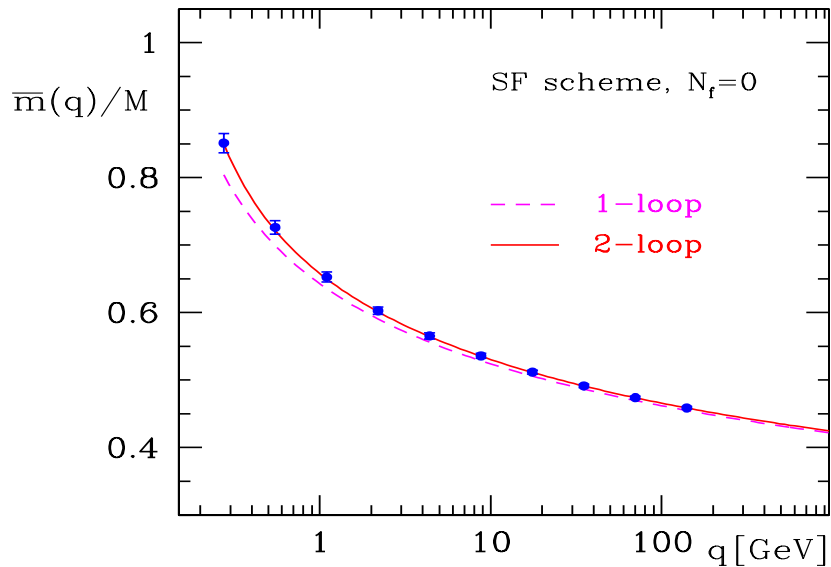


Fundamental Parameters of QCD

- running coupling $\alpha_s(E)$



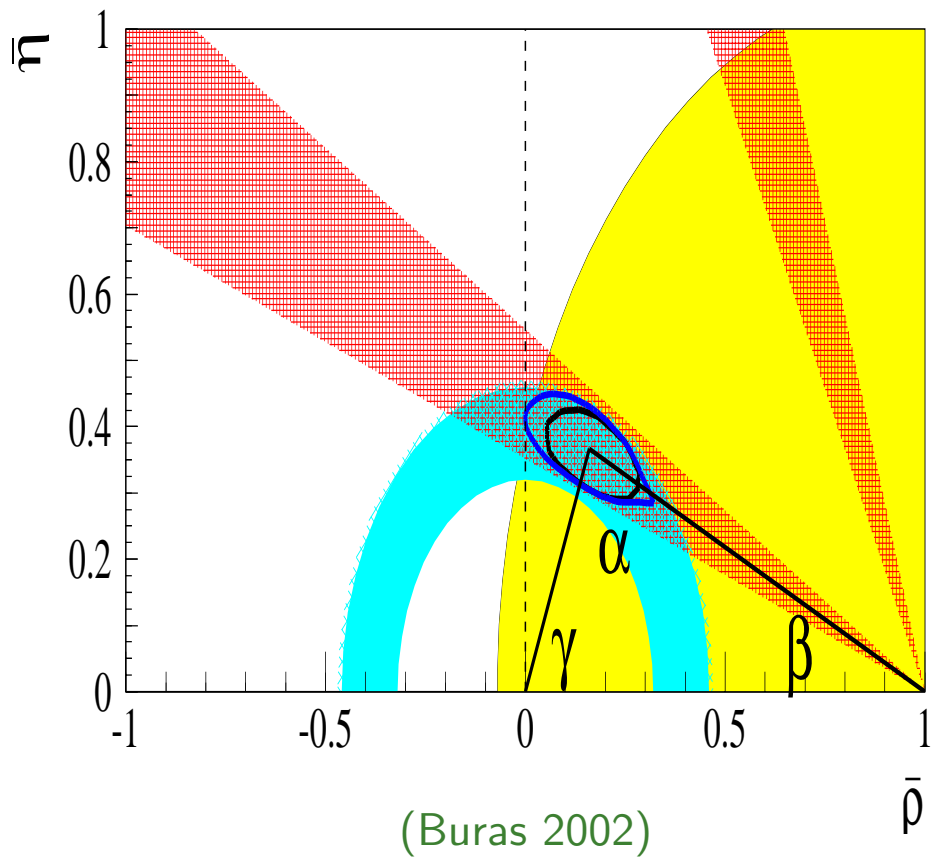
- running quark mass $m(q)$



non specialist introduction

R. Sommer and H. Wittig, physics/0204015

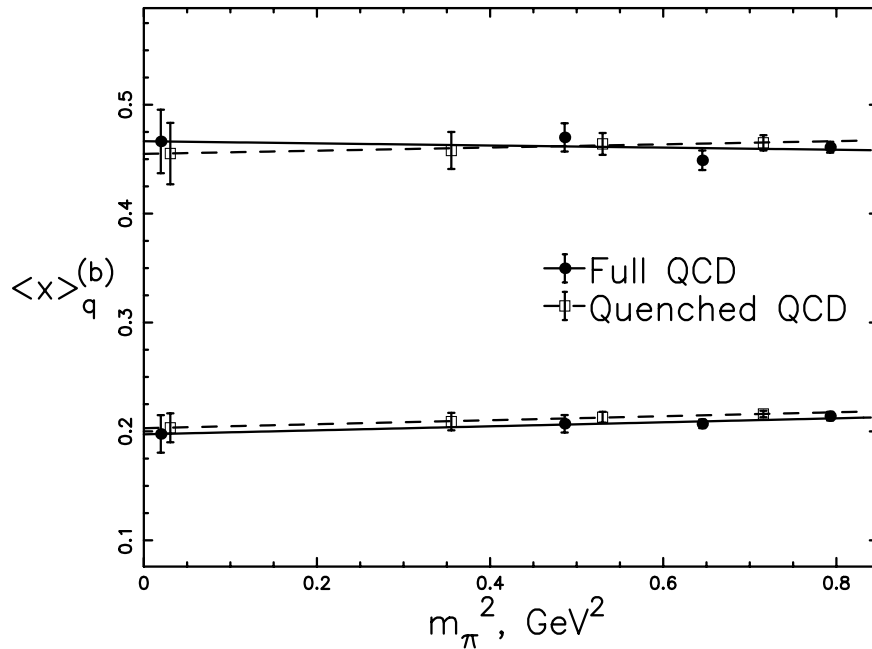
Unitarity triangle



→ sides of triangle constrained by results of lattice calculations

Moments of Parton distribution functions

example: lowest moment of twist-2, non-singlet operator in pion



LHPC, MIT-Wuppertal, QCDSF collaboration

quenched numbers:

Guagnelli, Jansen, Palombi, Petronzio, Shindler, Wetzorke

$$\langle x \rangle^{\text{experiment}} (\mu = 2.4 \text{ GeV}) = 0.23(2)$$

$$\langle x \rangle_{\overline{\text{MS}}}^{\text{quenched}} (\mu = 2.4 \text{ GeV}) = 0.30(3)$$

Cost of numerical simulations

Quenched fermions

expensive part: fermion (quark) propagator D^{-1}

$$\langle \bar{\psi}(x)\psi(y) \rangle \propto D^{-1}b$$

b is external source vector

⇒ need to solve

$DX = b$ with D a complex matrix that is

- high-dimensional $O(10^6) \otimes O(10^6)$
- sparse (diagonal and a few subdiagonals)

example: quenched $16^3 \cdot 32$ lattice:

≈ 100Mflops for one fermion matrix times vector multiplication

having a 50 Gflops (sustained) machine

⇒ about 10 hours for a physical result at one set of parameters

realistic lattices today: $32^3 \cdot 64$ or $48^3 \cdot 96$

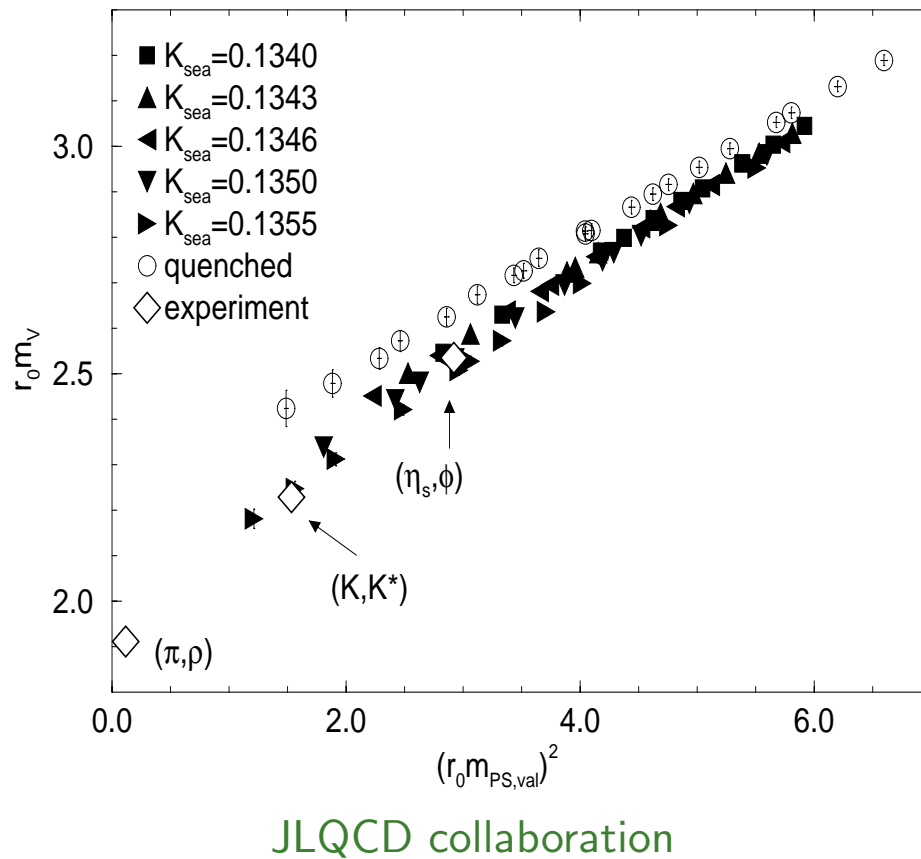
→ factor 10 to 100

Dynamical fermions

→ additional factor of $O(100)$

First results with dynamical fermions

example: vector meson spectrum



- effects of dynamical quarks visible
- systematics not under control yet
 - *continuum limit*
 - *finite size effects*
 - *ρ -meson decay*

Dynamical fermions

generate new configuration by integrating the field equations in a *Monte Carlo time* τ

$$\begin{aligned}\delta\pi/\delta\tau &= -\delta S/\delta U \\ \delta U/\delta\tau &= \pi\end{aligned}$$

fields: *gauge fields* U
conjugate momenta π

numerical integration \Rightarrow non-vanishing step size $\delta\tau$

experience: $\delta\tau N_\tau = 1$, $\delta\tau \approx 0.01$

discrete integration step size \Rightarrow integration error
 \Rightarrow global *accept/rejection step*

The difficulty:

$\delta_U S \propto [D^\dagger D]^{-1} \Phi$, Φ Gaussian random vector

since $N_\tau \approx 100$

\Rightarrow simulations are $O(100)$ more expensive than quenched approximation

Costs of dynamical fermion simulations

see panel discussion in Lattice2001, Berlin, 2001

Cost of 1000 dynamical $N_F = 2$ configurations

$$C_{\text{per}} = F_{\text{per}} \left(\frac{m_\pi}{m_\rho} \right)^{-z_\pi} \left(\frac{L}{a} \right)^{z_L} \left(\frac{r_0}{a} \right)^{z_a}$$

$$F_{\text{per}} = 6 \cdot 10^6 \text{ flops}$$

$$z_\pi = 6, \quad z_L = 5, \quad z_a = 2$$

given a 10 Teraflops computer,

lattice spacing $a = 0.1\text{fm}$, physical volume of 2fm^4

– for $m_\pi/m_\rho = 0.5 \rightarrow 70$ days

– for $m_\pi/m_\rho = 0.4 \rightarrow 270$ days

- do not simulate directly at the physical point
- combine simulation results with analytical methods to extrapolate
 \rightarrow *chiral perturbation theory*

question of overlap region of simulation and chiral perturbation theory major research topic is $m_\pi/m_\rho = 0.4$ sufficient?

From enemies to friends: Chiral symmetry on the lattice

chiral symmetry (exchange of massless left and right-handed fermions) very important to explain low-energy phenomena

in the continuum:

chiral symmetry expressed as $D_{\text{cont}}\gamma_5 + \gamma_5 D_{\text{cont}} = 0$

on the lattice: different anti-commutation relation

$$\begin{aligned}\gamma_5 D_{\text{latt}} + D_{\text{latt}} \gamma_5 &= 2a D_{\text{latt}} \gamma_5 D_{\text{latt}} \\ \rightarrow \gamma_5^{\text{cont}} \rightarrow \gamma_5^{\text{latt}} &= \gamma_5^{\text{cont}} (1 - a D_{\text{latt}})\end{aligned}$$

realizations of such a D_{latt}

- overlap operator (Neuberger)
- domain wall fermions (Kaplan, Shamir)
- fixed point action (Hasenfratz, Niedermayer, Wiese)

chiral invariant formulations of lattice QCD

↑ enjoy many properties of continuum theory

↑ can reach very small quark mass region

↓ are $O(100)$ more expensive than standard lattice fermions

Scalar condensate

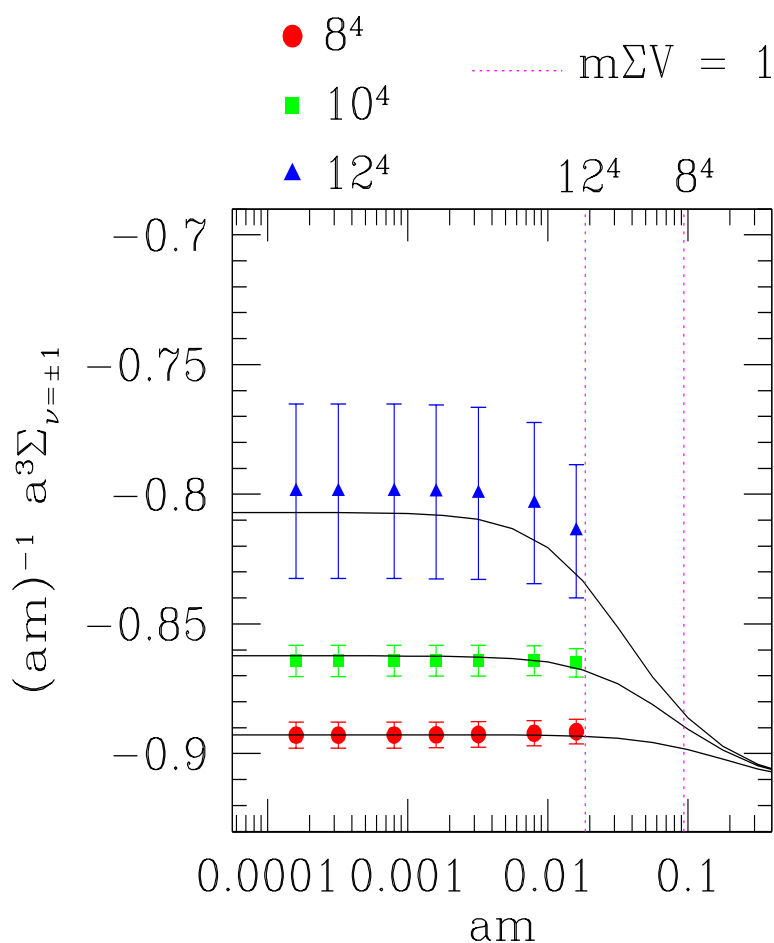
data points at 7 masses on 3 volumes

attempt a fit according to chiral perturbation theory

$$\Sigma_{\nu=\pm 1} = m\Sigma^2V + O((m\Sigma V)^2) + C/a^2$$

→ fixed topological sector $\nu = \pm 1$

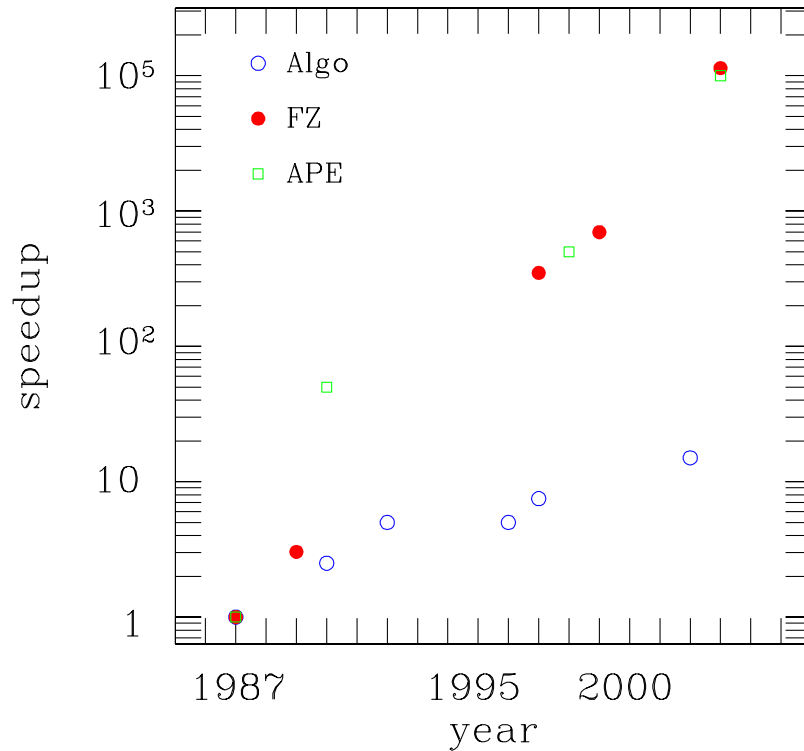
→ only two free parameters infinite volume
scalar condensate Σ and divergence C



(Hernández, Lellouch, Jansen, hep-lat/0203029)

⇒ find strong evidence for spontaneous chiral symmetry
breaking in QCD!

Algorithm and Machine Development



○ : algorithm development: factor 20

● : FZ CRAY at Research centre Jülich

□ : “Array Processor Experiment” (APE)

– machine development most important

– APE and CRAY equal performance
(following both Moore’s law)

– $\text{cost}(\text{CRAY}) \approx 10 \times \text{cost}(\text{APE})$

⇒ worthwhile to build machines

APE (*Europe*), QCDOC (*USA*), PC-Cluster

Japan

Computational Physics on Parallel Array Computer System

→ CPPACS

collaboration of lattice physicists from Tsukuba

+ industrial partner Hitachi



614 Gflops peak speed

128 Gbytes memory

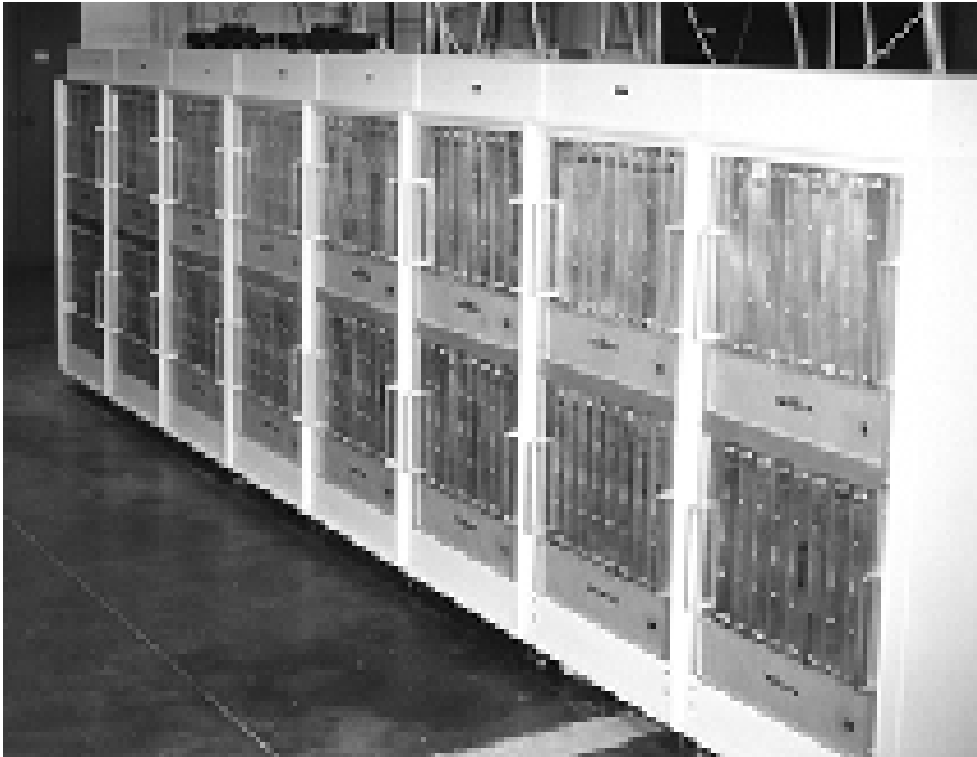
2048 Processing units

future development → ?

USA

QCD on digital Signal Processor System

→ QCDSP



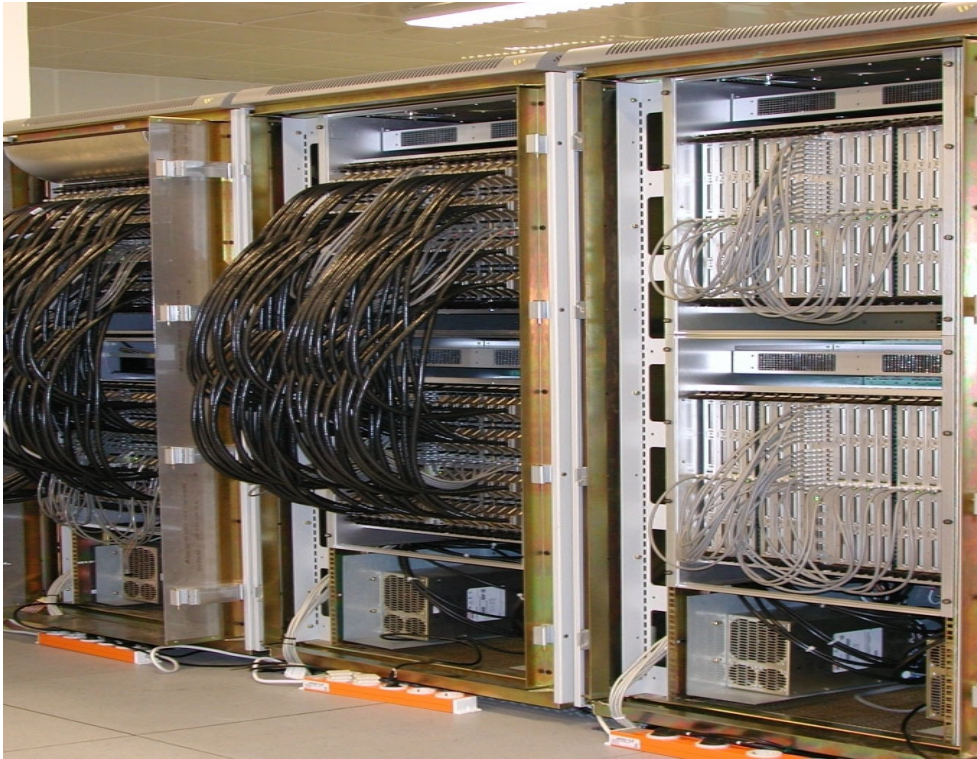
600 Gflops peak speed
50 Gbytes memory
12 288 Processing units

future development → QCDOC (QCD On Chip)
collaboration of lattice physicists from Columbia University, RIKEN,
BNL and UKQCD + industrial partner IBM

10 Tflops peak speed
40Gbytes on chip + $O(1)$ Tbytes external memory
 $O(10\ 000)$ Processing units
\$1/Mflops sustained performance

Europe

Array Processor Experiment → APE



APEmille installation in Zeuthen

550 Gflops peak speed

32 Gbytes memory

1024 Processing units

future development → apeNEXT

collaboration of lattice physicists from INFN, DESY
and University of Paris Sud

10 Tflops peak speed

1-4 Tbytes memory

O(6 000) Processing units

1€/Mflops sustained performance

Future needs

large machines \Rightarrow collaborative efforts

USA

Scientific Discovery through Advanced Computing
SciDAC

demand of about 40 Teraflops realized by

- QCDOC
- large PC-Cluster installations

evaluation of needs in Europe

- ECFA Report
Requirements for high performance computing for lattice QCD: report of the ECFA working panel
F. Jegerlehner et.al., CERN 2000-002, ECFA/00/200
- NuPECC Report
The NuPECC Working Group Computational Nuclear Physics
M. Baldo et.al., June 2000



→ special situation in Germany

Forum of German lattice physicists

+ *association of groups in Austria & Switzerland*

- common initiative of lattice physicists
Universities, GSI and NIC/DESY
 - rich and diverse spectrum of physics
 - *Fundamental parameters of QCD*
 - *Hadron spectrum*
 - *Structure functions*
 - *Physics of B-mesons*
 - *QCD thermodynamics*
 - *QCD at non-vanishing baryon density*
 - *Supersymmetry on the lattice*
 - *Algorithms*
 - *Chiral invariant lattice QCD*
- detailed definitions of milestones
 - ⇒ integrated need of 25 Teraflops (peak)
- LATFOR evaluation group
development of benchmark suite for evaluating platforms
apeNEXT, QCDOC, PC-Cluster, CRAY, Hitachi, IBM

The **John von Neumann-Institute of Computing (NIC)**

cooperation between **DESY** and **research centre Jülich**

- **NIC** shall provide supercomputer resources

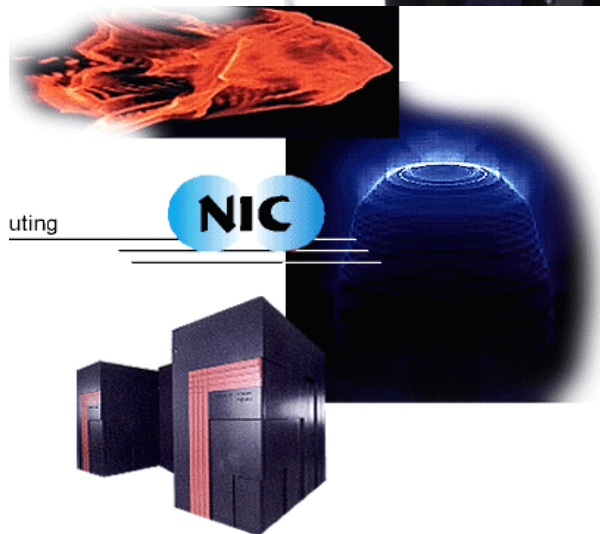
Zeuthen: centre of LGT

NIC research group

+Theory

+Universities **HU, FU**

+APE100, APEmille



Jülich: centre of general computational science

Supercomputers: **CRAY, IBM**

- **NIC** research groups
 - Many particle physics (P. Grassberger)*
 - Elementary particle physics (K.J.)*

Numerical simulations

Monte Carlo integration

compute partition function (Feynman's pathintegral)

$$Z = \int_{\text{fields}} e^{-S}$$

$$\int_{\text{fields}} = \prod \int \text{degrees of freedom}$$

with **degrees of freedom** = $O(10^6)$ - $O(10^8)$

→ *Monte Carlo method with importance sampling*
(Metropolis, heatbath, over-relaxation, cluster, ...)

propagator measurement

$$\langle \bar{\psi}(x)\psi(y) \rangle \propto D^{-1}b$$

b is external source vector

⇒ need to solve

$DX = b$ with D a complex matrix that is

- high-dimensional $O(10^6) \otimes O(10^6)$
- sparse (diagonal and a few subdiagonals)

- ★ modern lattice computations
- do not only want to have bigger computers
- work hard on algorithmic improvements
- incorporate theoretical progress:
 - continuum limit $a \rightarrow 0$
 - ← only acceleration of approach to the continuum limit
 - have developed *exact chiral symmetry on the lattice*: important theoretical (numerical?) concept
 - get rid of effects of finite physical boxlength L
 - ⇐ use the finite extend of the box
 - Finite Size Scaling technique*

on the machine side:

- race between **apeNEXT** and **QCDOC**
- question: *role of PC-clusters*

big challenge is present transition towards dynamical fermion simulations

- exciting → real physics
- powerful supercomputers are an essential ingredient
- combination with analytical methods are equally necessary